

## CFD INVESTIGATION OF FLOW OVER A BACKWARD-FACING STEP USING AN RNG $k$ - $\varepsilon$ TURBULENCE MODEL

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### ABSTRACT

Backward-facing step (BFS) is a benchmarked geometry for visualizing recirculation flow and validating turbulence models. Nowadays, numerical analysis with the CFD method has become more popular and has stimulated research involving CFD without avoiding the experimental method. In this paper, flow over a BFS was numerically investigated with an RNG  $k$ - $\varepsilon$  turbulence model to predict recirculating flow. BFS geometry refers to the geometry proposed by Kasagi & Matsunaga; it is three-dimensional, with inlet  $Re = 5.540$ . The paper aims to investigate the performance of the RNG  $k$ - $\varepsilon$  turbulence model over a BFS. Two important parameters were analyzed: the performance of the RNG  $k$ - $\varepsilon$  on the recirculation zone and on the reattachment length. Recirculation flow is presented by the x-velocity for  $Y = 17.4$  mm and  $Y = 34.9$  mm. In these Y-section, the RNG  $k$ - $\varepsilon$  is compared to the STD  $k$ - $\varepsilon$  and both models show the recirculation flow occurred from  $X = 0$  mm to about  $X = 200$  mm. The following results were obtained. The RNG  $k$ - $\varepsilon$  predicted a slightly higher x-velocity component than that predicted by the STD  $k$ - $\varepsilon$ . This result shows that the RNG  $k$ - $\varepsilon$  turbulence model is suitable for predicting recirculation flow on the BFS. The reattachment length was measured by non-dimensional  $X/h$  to the x-velocity component with the RNG  $k$ - $\varepsilon$  turbulence model. The analyzed data were taken from  $X/h = 4.5$  to  $X/h = 10$ , on the x-velocity component from  $Y = 17.4$  mm. The reattachment point was achieved at  $X/h = 7.22$ , close to that achieved by Kasagi & Matsunaga of  $X/h = 6.51$ .

*Keywords:* Backward-facing step; CFD; Reattachment point; Recirculation flow; RNG  $k$ - $\varepsilon$  turbulence model

### 1. INTRODUCTION

Backward-facing step (BFS) is the one of the most powerful geometries for visualizing flow, validating the performance of turbulence model on recirculating flow (Thangam & Speziale, 1991; Thangam, 1991). Generally, there are two main specific flows in a BFS considered as a benchmarking geometry: the recirculating flow after the expansion zone and the reattachment point reaching near to the outlet zone. The recirculated and swirling flow occurs in many engineering applications well-presented by BFS geometry. This type of flow can be useful or harmful, depending on the application; examples include recirculation flow in electronic devices; recirculation flow in aerodynamics fields; flow around buildings in architectural applications; flow in combustion chambers; and the disadvantage of swirling flow at pipe bends (Mouza et al., 2005; Rouizi et al., 2009; Gautier & Aider, 2014; Ramšak, 2015; Saha & Nandi, 2017; Selimefendigil & Öztöp 2017). There remain many turbulent flow phenomena over a BFS which

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are yet to be explored (Kasagi & Matsunaga 1995; Avancha, 2002).

There are several geometrical aspects to BFS, most of which are non-dimensional parameters, such as step height ( $h$ ), upstream height ( $H$ ), the expansion ratio and total length ( $L$ ). Several papers have investigated turbulent flow over a BFS, with specific cases examined experimentally and/or numerically. A review of these geometry parameters was previously made by (Darmawan, 2016). Experimental method of flow over BFS geometry as done by (Kasagi & Matsunaga, 1995; Gautier & Aider, 2013) and many others involving advanced measurement techniques and recording facilities, which are very costly (Gautier & Aider 2014; Kasagi & Matsunaga 1995). Moreover, flexibility in varying the geometry design, lower cost, and faster and better visualization have made the CFD method more popular over the years, two dimensionally and three dimensionally (Thangam, 1991; Avancha & Pletcher, 2002; Nie & Armaly, 2002; Kanna & Das, 2006; Rouizi et al., 2009; Ramšak, 2015). The growth of commercial CFD codes in the market is also increasing the number of CFD applications, with choices of characteristics, accuracy. Turbulence model choice plays an important role in producing the acceptable results, needing a physical flow and mathematical knowledge to perform CFD simulation (Ramdhan et al., 2016). An appropriate turbulence model is needed in order to represent the flow with commercial CFD codes. For example, in 2007 Anwar-ul-Haque et al. assessed the performance of turbulence models on backward-facing step applications (Haque et al., 2007).

Many turbulence models are available, with very wide range characteristics, from zero equation to DNS (Direct Numerical Simulation). The two-equation turbulence models (RANS-based) are often used in research regarding accuracy and computational resources compared to more complicated models (Thangam, 1991). The STD  $k-\varepsilon$  turbulence model is the most used model despite its weakness in presenting swirl-dominated flow (Launder & Spalding, 1974; Lakshminarayana, 1996; Marshall & Bakker, 2003). Another turbulence model based on the two-equation model is the RNG  $k-\varepsilon$ , with improvement in swirling and recirculating flow by renormalizing small scale eddies compared to the STD  $k-\varepsilon$  (Yakhot & Orszag, 1986; Thangam & Speziale, 1991; Versteeg & Malalasekera, 2007; Budiarto et al., 2013; Darmawan et al., 2013). Therefore, the RNG  $k-\varepsilon$  may be appropriate for predicting the recirculation flow and reattachment point of the flow on the BFS geometry.

However, the performance of the RNG  $k-\varepsilon$  turbulence model needs to be compared with the STD  $k-\varepsilon$  model, as the most used one, with faster computation and lower computational resources. The results then compared with the experimental results obtained by (Kasagi & Matsunaga 1995). The BFS geometry used here also refers to the geometry proposed by Kasagi and Matsunaga, but the effects of the boundary layer are not considered here. Therefore, this paper aims to investigate the performance of RNG  $k-\varepsilon$  turbulence model over a BFS. The results may be used as a reference and may be applied in future research involving flow in BFS geometry.

## 2. METHODS

### 2.1. Geometry Model

The BFS model consists of a number of important dimension parameters, most of which can be treated as non-dimensional parameters: *step height* ( $h$ ), *upstream height* ( $H$ ), expansion, and total length ( $L$ ). The expansion ratio is defined as  $((H+h)/H)$ . The geometry model used here refers to a model proposed by (Kasagi & Matsunaga, 1995). That employed an experimental method on the BFS with a 3D-PTV, with water as the working fluid. Table 1 shows the geometry specification.

Table 1 BFS geometry specification

Backward-facing step geometry			
Step height, $h$ (mm)	Upstream height, $h$ (mm)	Expansion ratio	Total Length, $L$ (mm)
41	81	1.5	$82.5H$
Width (mm)	20 $h$		
Working fluid	:	Water	

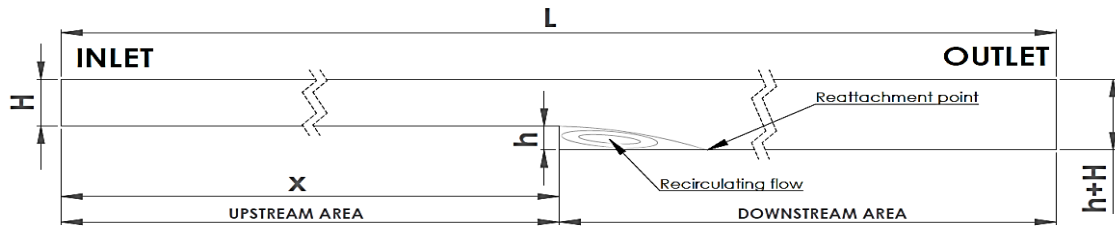


Figure 1 BFS Geometry

## 2.2. CFD Model

The computational grid that is used for CFD numerical simulation with Autodesk CFD 2017 commercial codes (academic version). Water was used as the working fluid, with the temperature for all four faces of the geometry set at 298K. According to (Kasagi & Matsunaga, 1995), the inlet Reynolds number is set to 5.540, with water as the working fluid. Computation is performed by using a three-dimensional unstructured mesh. Since there are no external forces involved, turbulent intensity is assumed to be 10%. The flow is assumed to be steady-state and the wall is assumed to have zero-roughness. This is also confirmed by the experimental study which is part of the numerical study, which used acrylic as the main material of the BFS geometry (Darmawan, 2017). Furthermore, the surface roughness of the acrylic is very low ( $0.07 \mu\text{m}$ ), therefore the wall can be assumed to have zero roughness (Rao et al., 2015).

Mesh dependency analysis was conducted, subject to the mass flow rate of two mesh numbers: 3,055 nodes and 10,426 nodes. Table 2 shows that the nodes amount almost not affect the prediction of the mass flow rate. Considering that the computational time needed was only slightly different, the computational grid with 10,426 nodes was chosen, as shown by Figure 2. Figure 3 shows the closer computational mesh from Figure 2 at the expansion zone. The computational boundaries used in both the mesh dependency analysis and actual analysis are follows: water as working fluid with Reynolds number ( $Re$ ): 5.540 and temperature of 298 K; zero roughness wall; and pressure at the outlet set to 0.1 MPa.

Table 2. Mesh dependency analysis

	RNG $k-\varepsilon$ turbulence model			
	Nodes: 3,055		Nodes: 10,426	
	Inlet	Outlet	Inlet	Outlet
Mass flow rate (kg/s)	2.293	-2.288	2.222	-2.288

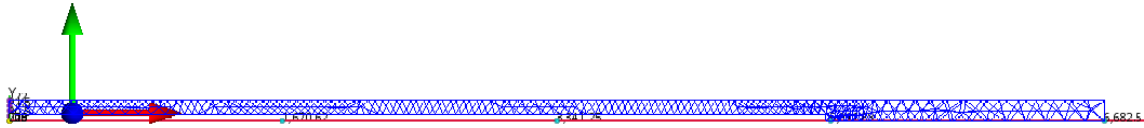


Figure 2 Computational grid of the BFS

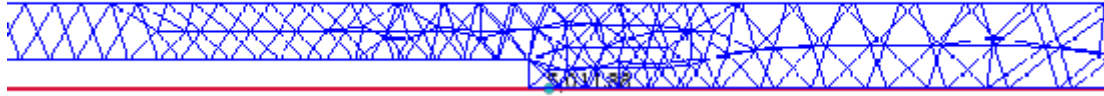


Figure 3 Computational grid at the expansion zone

### Governing Equation

As the most common model, the STD  $k$ - $\varepsilon$  turbulence model developed by Launder & Spalding is a two-equation (RANS-based) turbulence model, consisting of  $k$  and  $\varepsilon$  transport equations for handling turbulence kinetic energy and dissipation respectively (Launder & Spalding 1974)

$k$  transport equation

$$\frac{Dk}{Dt} = \frac{1}{\rho} \frac{\partial}{\partial x_k} \left[ \frac{\mu_T}{\sigma_k} \frac{\partial k}{\partial x_k} \right] + \frac{\mu_T}{\rho} \left( \frac{\partial U_i}{\partial x_k} + \frac{\partial U_k}{\partial x_i} \right) \frac{\partial U_i}{\partial x_k} - \varepsilon \quad (1)$$

$\varepsilon$  transport equation

$$\frac{D\varepsilon}{Dt} = \frac{1}{\rho} \frac{\partial}{\partial x_k} \left[ \frac{\mu_T}{\sigma_\varepsilon} \frac{\partial \varepsilon}{\partial x_k} \right] + \frac{C_1 \mu_T}{\rho} \frac{\varepsilon}{k} \left( \frac{\partial U_i}{\partial x_k} + \frac{\partial U_k}{\partial x_i} \right) \frac{\partial U_i}{\partial x_k} - C_2 \frac{\varepsilon^2}{\kappa} \quad (2)$$

Model constants

$C_\mu$	$C_1$	$C_2$	$\sigma_k$	$\sigma_\varepsilon$
0.09	1.44	1.92	1.0	1.3

The RNG  $k$ - $\varepsilon$  turbulence model was developed from the STD  $k$ - $\varepsilon$  turbulence model, in which the small-scale eddies are normalized (Renormalization Group) was developed by (Yakhot & Orszag, 1986). Mean flow property and turbulent effects are averaged (Versteeg & Malalasekera, 2007), (TM-107, 1993), (Munson et al., 2009). Reynolds stress governs the eddy viscosity concept, with reference to the Boussinesq equation (Blazek, 2005). With the application of renormalization theory, small scale eddies are computed and averaged separately and then solved together with the large-scale eddies. With this concept, the computational resources needed by the RNG turbulence model are similar to the STD  $k$ - $\varepsilon$ . Recirculation and swirling flows are turbulent flows involving wide range interaction between small- and large-scale eddies. This is the reason why the RNG  $k$ - $\varepsilon$  turbulence model can produce better results than the STD  $k$ - $\varepsilon$  turbulence model in handling recirculation flow. Like the other two-equation turbulence models, the RNG  $k$ - $\varepsilon$  turbulence model consists of  $k$  and  $\varepsilon$  transport equations for handling turbulence kinetic energy and dissipation respectively.

$k$  transport equation

$$\frac{\partial K}{\partial t} + (\bar{v} \cdot \nabla) K = \frac{\nu_T}{2} \left( \frac{\partial \bar{v}_i}{\partial x_j} + \frac{\partial \bar{v}_j}{\partial x_i} \right)^2 - \bar{\varepsilon} + \frac{\partial}{\partial x_i} \alpha_K \nu_T \frac{\partial K}{\partial x_i} \quad (3)$$

$$\frac{\partial K}{\partial t} + (\bar{v} \cdot \nabla) K = 0.3 (2S_{ij}^2)^{1/2} - \bar{\varepsilon} + \frac{\partial}{\partial x_i} \alpha_K \nu_T \frac{\partial K}{\partial x_i} \quad (4)$$

$$\frac{\partial K}{\partial t} + (\bar{v} \cdot \nabla) K = P - \bar{\varepsilon} + \frac{\partial}{\partial x_i} \alpha_K \nu_T \frac{\partial K}{\partial x_i} \quad (5)$$

$\varepsilon$  transport equation

$$\frac{D\bar{\epsilon}}{Dt} = \bar{P} - 1.7215 \frac{\bar{\epsilon}^2}{K} + \frac{\partial}{\partial x_i} \alpha_\epsilon \nu \frac{\partial \bar{\epsilon}}{\partial x_i} \tag{6}$$

$$\frac{D\bar{\epsilon}}{Dt} = -1.063 \frac{\bar{\epsilon}}{K} \tau_{ij} \frac{\partial \bar{v}_i}{\partial x_j} - 1.7215 \frac{\bar{\epsilon}^2}{K} + \frac{\partial}{\partial x_i} \alpha_\epsilon \nu \frac{\partial \bar{\epsilon}}{\partial x_i} \tag{7}$$

General form for  $\epsilon$  transport eq.

$$\frac{D\bar{\epsilon}}{Dt} = -C_{\epsilon 1} \frac{\bar{\epsilon}}{K} \tau_{ij} \frac{\partial \bar{v}_i}{\partial x_j} - C_{\epsilon 2} \frac{\bar{\epsilon}^2}{K} + \frac{\partial}{\partial x_i} \alpha_\epsilon \nu \frac{\partial \bar{\epsilon}}{\partial x_i} \tag{8}$$

Model Constants:

$C_\mu$	$C_{\epsilon 1}$	$C_{\epsilon 2}$	$Pr_t$	$C_\kappa$	$\kappa$	$Ba$	$S_3$
0.0845	1.063	1.7215	0.7179	1.617	0.372	1.161	0.478

where  $C_\mu$  is the total viscosity constant,  $C_{\epsilon 1}$  is the turbulent production constant on dissipation transport,  $C_{\epsilon 2}$  is the dissipation constant on dissipation transport,  $\kappa$  is the von Kármán constant,  $C_\kappa$  is the constant,  $Ba$  is the Batchelor constant,  $S_3$  is the matrix skewness factor.

### 3. RESULTS

CFD simulation was conducted with Autodesk CFD 2017 commercial codes. Figure 4 to Figure 7 show the CFD simulation results with STD  $k-\epsilon$  and RNG  $k-\epsilon$ . The two specific flow parameters on the BFS analysed were recirculation flow and reattachment length, represented by Figure 5 and Figure 7. Based on these reasons, analysis subjected to the  $v_x$  velocity component (inlet to outlet direction). The recirculation flow is characterized by the intensity of the reverse direction of the  $v_x$  velocity component, while the reattachment length is measured by the distance from the expansion zone start to the point at which the  $v_x$  velocity component is reversing ( $X$ ) compared to step height ( $h$ ) = ( $X/h$ ).

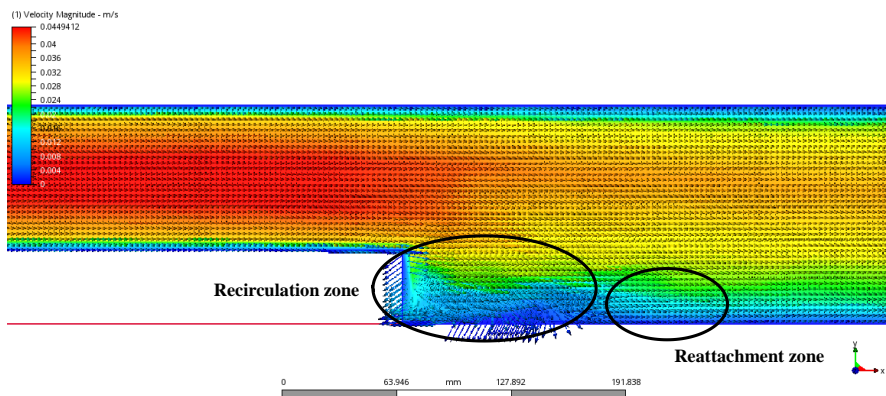


Figure 4 STD  $k-\epsilon$ ,  $v_x$  velocity component

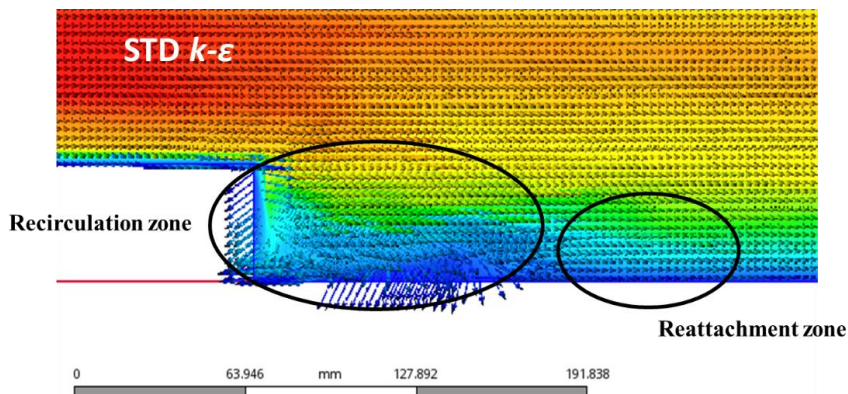
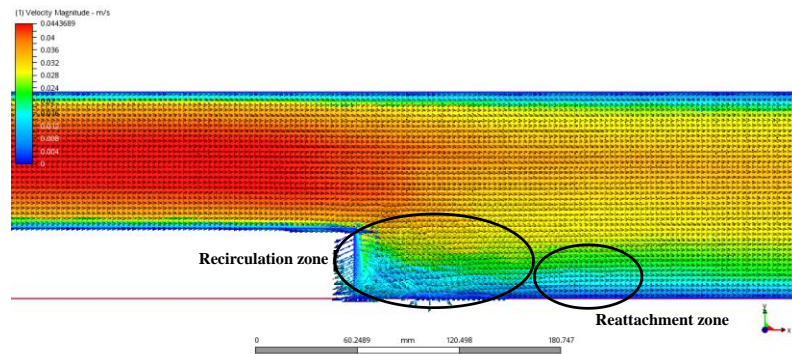
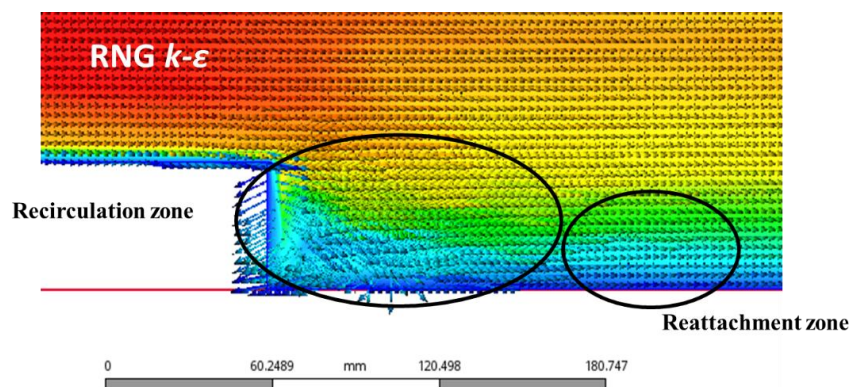


Figure 5 Recirculation and reattachment zone computed with the STD  $k-\epsilon$  turbulence model

Figure 6 RNG  $k-\epsilon$ ,  $v_x$  velocity componentFigure 7 Recirculation and reattachment zone computed with the RNG  $k-\epsilon$  turbulence model

### 3.1. STD $k-\epsilon$ and RNG $k-\epsilon$ Turbulence Model Comparison

Figure 4 and Figure 6 show the x-y plane with the  $v_x$  velocity component generated by the STD  $k-\epsilon$  and RNG  $k-\epsilon$  turbulence models respectively. The plane is located at the midpoint of the BFS width. Since the width of the BFS is large enough and the upstream length is long enough, this plane can be considered to represent the general flow as the flow is fully developed so can be assumed to be a two-dimensional flow. With  $Re = 5,540$ , both models show similar results on the  $v_x$  velocity component on the upstream zone (inlet zone to the expansion zone). The different shown by the recirculation zone and reattachment length. The recirculation zone predicted by the RNG  $k-\epsilon$  turbulence model tends to generate more fluctuation between small scale and large scale eddies, with a higher velocity than that generated by the STD  $k-\epsilon$  turbulence model. The reattachment zone which reached lower on the RNG  $k-\epsilon$  might be caused by the fluctuation predicted by RNG  $k-\epsilon$  turbulence model. This phenomenon is in line with the principle of the turbulent flow which resulting an efficient mixing between smaller and larger eddies, cascade energy is decreasing on STD  $k-\epsilon$  turbulence model (Tennekes & Lumley, 1972). More detailed results are given in the following section, with the graphical forms.

## 4. DISCUSSION

### 4.1. Recirculation Zone

Recirculation flow occurs in the expansion zone, where the fluid flows from the upstream to the downstream zone. Figure 5 and Figure 7 show more detailed analysis made in this area by comparing the results generated by the STD  $k-\epsilon$  and RNG  $k-\epsilon$  turbulence models on a graph presented by  $v_x$  velocity component in the near-wall zone. A higher fluctuation velocity profile is assumed to give a better prediction of the recirculation zone. Figure 8 gives a clearer

understanding of the datum position used in the graphical analysis, with the datum point set on the left-hand corner side of the expanded area.

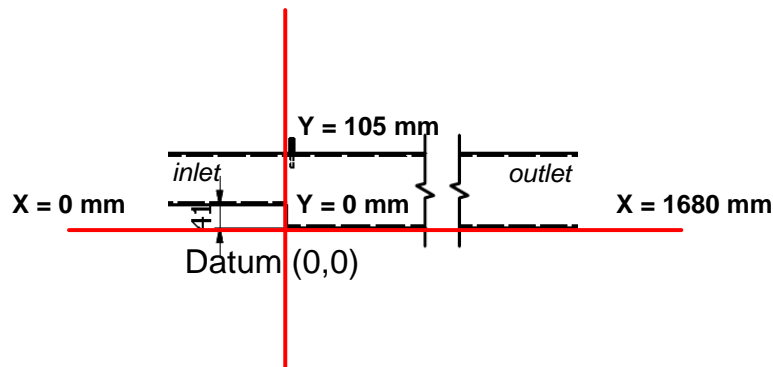


Figure 8 Simplified BFS geometry with datum point

The x-velocity component along the i-direction is presented graphically, starting from the datum point to the outlet side. Two Y-points are considered. Figure 9 and Figure 10 show the velocity profiles at  $Y = 17.4$  mm and  $Y = 34.9$  mm respectively. Both models show that the recirculation flow occurred from  $X = 0$  mm to about  $X = 200$  mm, as shown in Figure 4 and Figure 6. The main reason the STD  $k-\varepsilon$  turbulence model is used is because of its popularity as a fluid computational method. The velocity profile trend is similar for both models, but a higher x-velocity component is predicted by the RNG  $k-\varepsilon$  turbulence model than by the STD  $k-\varepsilon$  turbulence model, meaning that interaction between small- and large-scale eddies is higher, so the velocity is greater. Cascade energy is dominant in this condition, since the recirculation flow is closely related to the interaction between the scales of eddy. The additional terms of the RNG  $k-\varepsilon$  for both  $k$  and  $\varepsilon$  transport equations undoubtedly make the computational time with this model longer than that of the STD  $k-\varepsilon$ . Considering the cascade energy concept which larger velocity production represented by turbulent production term in Equation 3, the RNG  $k-\varepsilon$  turbulence model is more suitable for this type of flow.

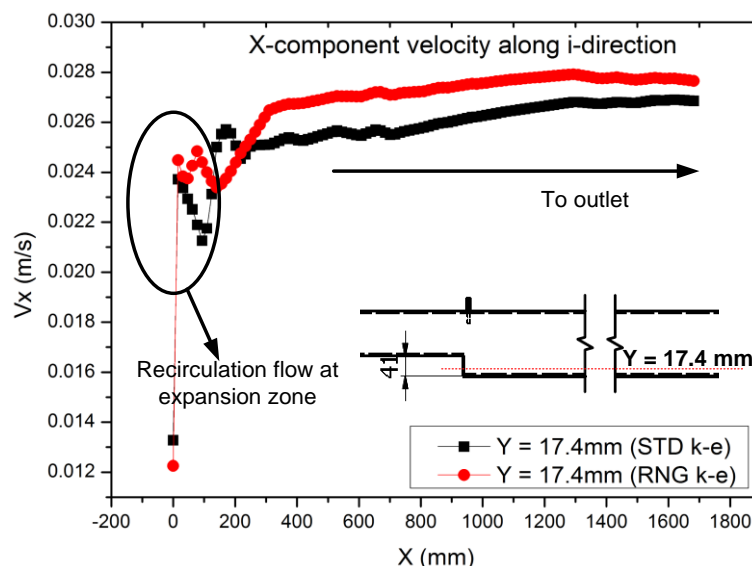


Figure 9 X-component velocity along i-direction,  $Y = 17.4$  mm

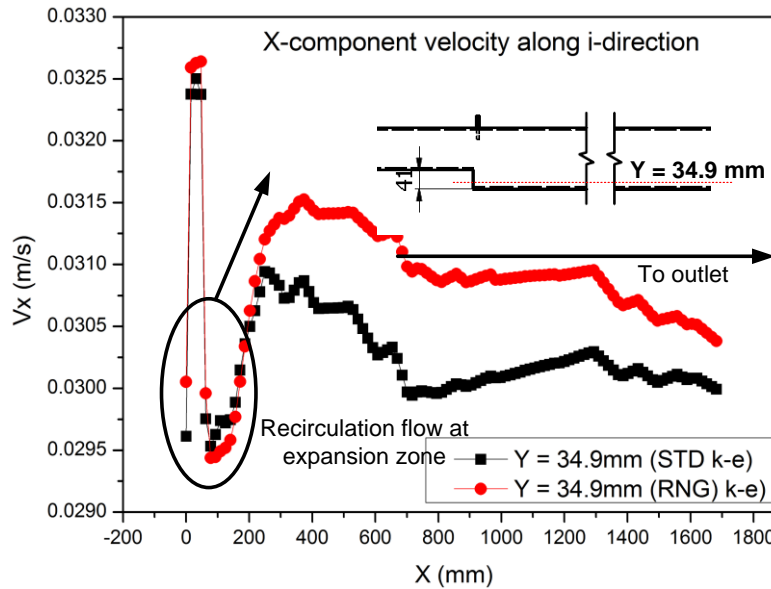


Figure 10 X-component velocity along i-direction, Y = 34.9 mm

#### 4.2. Reattachment Length

It is well-known that reattachment length is one of the most important flow parameters in BFS geometry. Hence, the reattachment point can be defined as the location at which the x-velocity component increases after recirculating and decreases to the outlet side smoothly. According to this definition, the reattachment length is measured by non-dimensional  $X/h$  to the x-velocity component. Analyzed data were taken from  $X/h = 4.5$  to  $X/h = 10$ , on the x-velocity component from  $Y = 17.4$  mm, as can be seen in Figure 11. With the RNG  $k-\epsilon$  turbulence model, the reattachment point is achieved at  $X/h = 7.22$ . This approach is different to the definition of Kasagi & Matsunaga, who define reattachment length as the location where the time fraction of forward flow is equal to 0.5 towards the wall (Kasagi & Matsunaga, 1995). However, the result is not far from the  $X/h = 6.51$  achieved by Kasagi & Matsunaga. The velocity profile tends to decline steadily after  $X/h = 10$ .

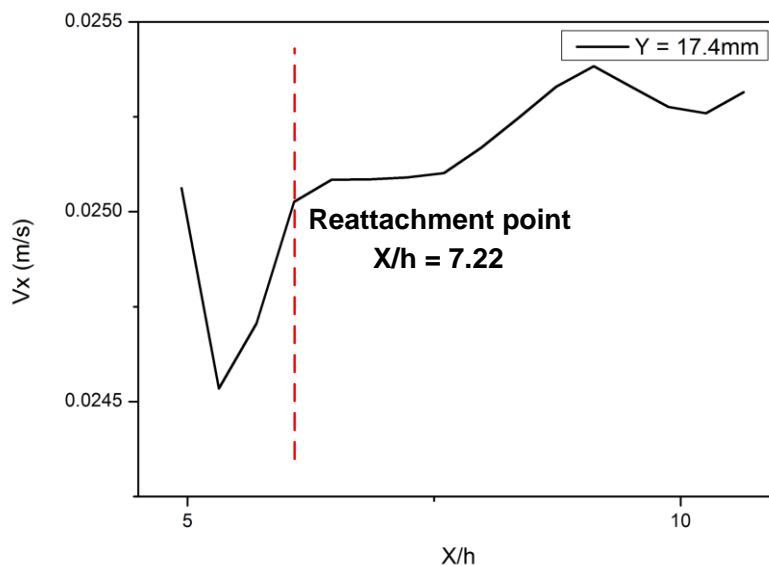


Figure 11 Reattachment length



## 5. CONCLUSION

A numerical study of the flow over a backward-facing step was conducted with  $Re = 5.540$  with STD  $k\varepsilon$  and RNG  $k-\varepsilon$  turbulence models. The conclusions are as follows: (1) The recirculation flow predicted by the RNG  $k-\varepsilon$  turbulence model is higher than that predicted by STD  $k-\varepsilon$  turbulence model in qualitative terms. This was measured by the x-velocity component along the x-direction, which shows that the RNG  $k-\varepsilon$  is better for use in such a flow; and (2) With the RNG  $k-\varepsilon$  turbulence model, at  $Y = 17.4$  mm the reattachment point was achieved at  $X/h = 7.22$ .

## 6. ACKNOWLEDGEMENT

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